

Original

P4820 Assignment IV

Due, April 6 (negotiable!)

- 1) A solid sphere of radius a is immersed in a vat of fluid at temperature T_0 . Heat is conducted into the sphere according to equation

$$\frac{\partial T}{\partial t} = \frac{kA}{mc} \nabla^2 T = D \nabla^2 T$$

where $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$

If the temperature of the boundary is fixed at T_0 , and the initial temperature of the sphere is T_1 , find the temperature within the sphere as a function of time. Plot the temperature $(T - T_0)/(T_1 - T_0)$ as a function of radius r for $Dt/a^2 = 1/20, 1/10, \text{ and } 1/5$.

Hint: The equation in r is "easier" to deal with if you make the change:

$$\frac{\partial}{\partial r} \left(r^2 \frac{\partial R}{\partial r} \right) = r \frac{\partial^2}{\partial r^2} (rR).$$

- 2) Consider a sphere of radius a for which the top half is set to voltage V and the bottom half is set to voltage 0 (grounded). Find the potential for this sphere for values of $r > a$ (ie. the potential on the outside of the sphere). You will recall that we did this problem solving for the potential on the inside of the sphere in class. Create a plot of the potential (contour?) to demonstrate that your solution makes sense, ie. voltages are appropriate at the surface of the sphere and they go to zero at infinity.
- 3) A cylinder of height h and radius a has the top and bottom grounded. The potential on the wall at $\rho = a$ is V_0 . Find the potential inside the cylinder. Plot $\Phi(\rho, z = h/2)/V_0$ for $0 < \rho < a$ and $\Phi(\rho = a/2, z)/V_0$, for $0 < z < h$.

1) Assignment 4 Answer

$$\frac{\partial T}{\partial t} = D \nabla^2 T$$

First use separation to get t out

$$T = z(R)$$

∇^2 in spherical -

$$R \frac{\partial z}{\partial t} = D \cdot \left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial rR}{\partial r} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial}{\partial \theta} \left(\sin^2 \theta \frac{\partial R^2}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 R^2}{\partial \phi^2} \right]$$

$$R \frac{\partial z}{\partial t} = D \cdot \frac{z}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial R}{\partial r} + \underbrace{0 + 0}$$

no angular dependence

$$\frac{1}{z} \frac{\partial z}{\partial t} = \frac{D}{R} \cdot \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial R}{\partial r} = k^2 \leftarrow \text{separation constant}$$

$$= \frac{D}{rR} \cdot \frac{1}{r^2} \cdot r^2 \frac{\partial^2 rR}{\partial r^2}$$

$$\frac{1}{z} \frac{\partial z}{\partial t} = \frac{D}{rR} \frac{\partial^2 rR}{\partial r^2} = k^2$$

For time

$$\frac{1}{\tau} \frac{\partial^2}{\partial t^2} = k^2$$

$$\ln \tau = k^2 t$$

$$\tau = e^{-k^2 t}$$

For range

$$\frac{D}{vR} \frac{\partial^2 vR}{\partial r^2} = -k^2$$

otherwise could not match initial condition!

$$\frac{\partial^2 vR}{\partial r^2} = \frac{-k^2}{D} (vR)$$

$$vR = A \sin \frac{kr}{D^{1/2}} + B \cos \frac{kr}{D^{1/2}}$$

..... This problem is eased if I form in terms of θ that $T_0 + \theta = \pi$ so that as $t \rightarrow \infty \theta \rightarrow T_0$

For time problem

Form problem as

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$$T - T_0 = \tau \cdot R$$

$$T = \tau \cdot R + T_0$$

and as $t \rightarrow \infty$ $\tau R \rightarrow 0$

also $T = T_0$ at $r = a$

also for $r \rightarrow a$ $T = T_0$

$$T - T_0 = e^{-k^2 t} \cdot \frac{1}{r} \left[A \sin \frac{k}{D^{1/2}} r + B \cos \frac{k}{D^{1/2}} r \right]$$

But... For r , need $R(r)|_{r=a} = 0$

set $B = 0$ and require $\frac{k}{D^{1/2}} a = n\pi$

$$k = \frac{n\pi D^{1/2}}{a}$$

general solution is

$$T - T_0 = e^{-\left(\frac{n\pi}{a}\right)^2 D t} \cdot A \frac{1}{r} \sin \frac{n\pi D^{1/2}}{a} r$$

$$\left. \frac{\cos ar}{r} \right|_{r=0} = \infty$$

$$\left. \frac{\sin ar}{r} \right|_{r=0} = 1$$

← B solutions are unbounded at $r=0$!

$$T - T_0 = A_n e^{-\left(\frac{n\pi}{a}\right)^2 Dt} \cdot \frac{1}{r} \sin \frac{n\pi}{a} r$$

at $t=0$, $T - T_0 = T_1 \quad \forall r$
 $T = T_1$

$$\sum_{n=0}^{\infty} A_n \frac{1}{r} \sin \frac{n\pi}{a} r = (T_1 - T_0) \quad r \leq a$$

$$\sum_{n=0}^{\infty} \int_0^a A_n \sin \frac{n\pi}{a} r \cdot \sin \frac{n\pi}{a} r \, da = \int_0^a (T_1 - T_0) r \sin \frac{n\pi}{a} r \, dr$$

$$\frac{A_n \cdot a}{2} = \int_0^a (T_1 - T_0) r \sin \frac{n\pi}{a} r \, dr$$

$$(T_1 - T_0) \left[r \frac{a}{n\pi} \cos \frac{n\pi}{a} r \right]_0^a$$

$$- (T_1 - T_0) \int_0^a \frac{-a}{n\pi} \cos \frac{n\pi}{a} r \, dr$$

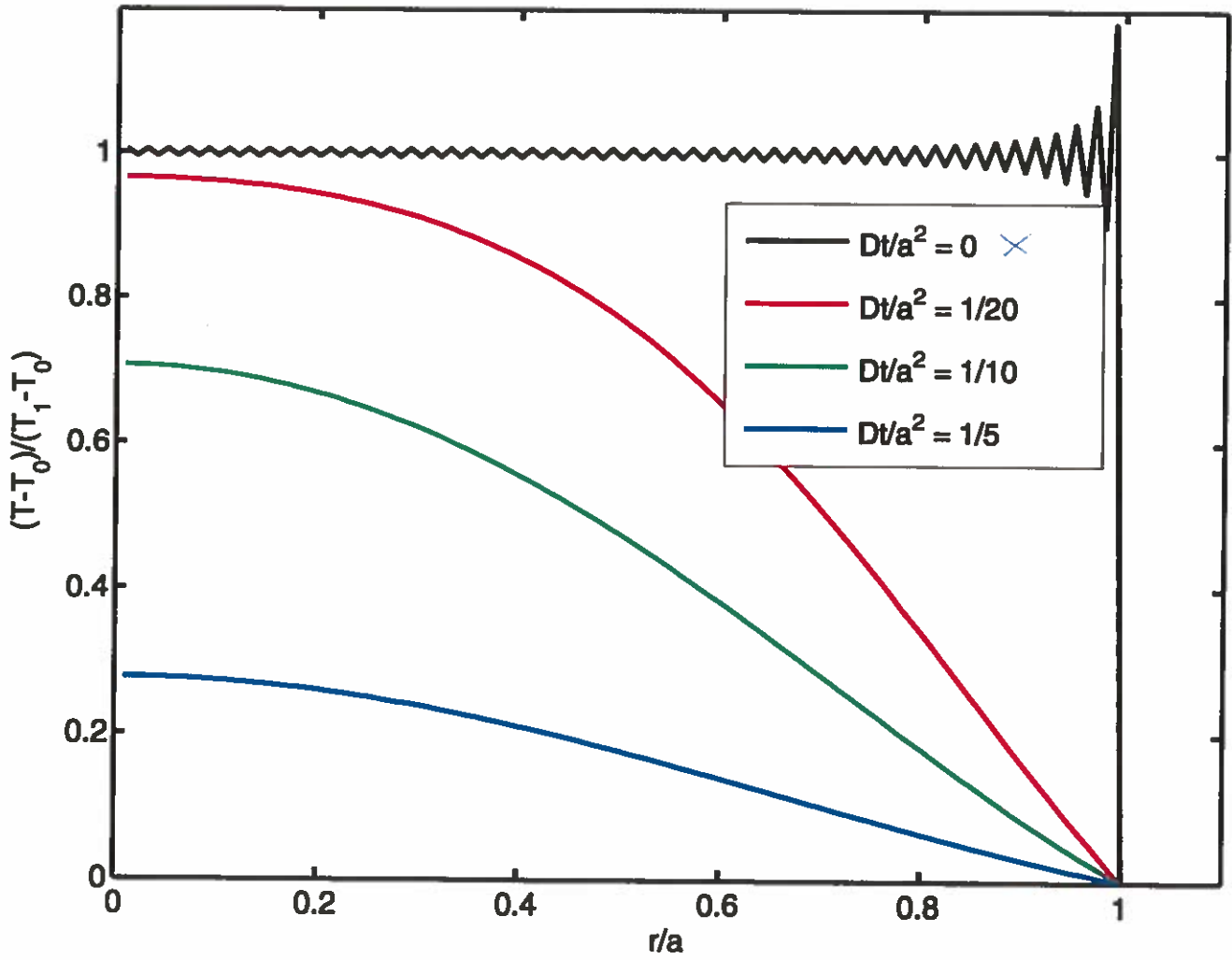
$$= - (T_1 - T_0) a^2 / n\pi \cos n\pi + (T_1 - T_0) \frac{a}{n\pi} \frac{a}{n\pi} \sin \frac{n\pi}{a} r \Big|_0^a$$

$$- (T_1 - T_0) \frac{a^2}{n\pi} (-1)^n + (T_1 - T_0) \frac{a^2}{n^2 \pi^2} \sin n\pi$$

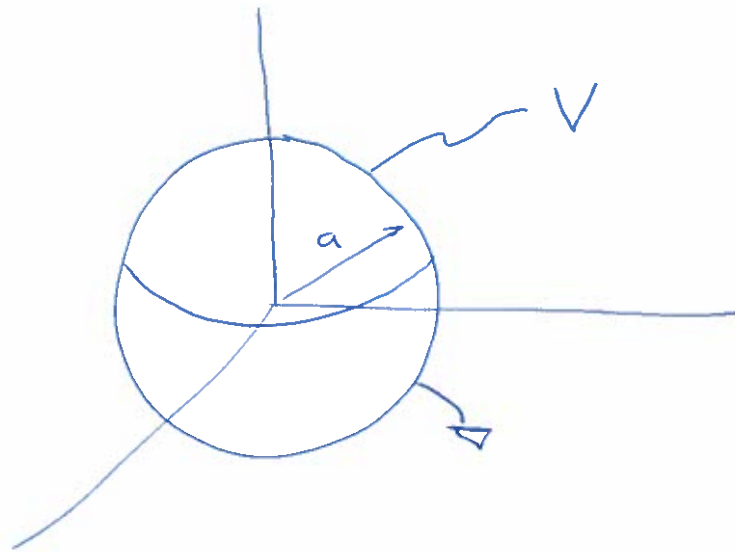
$$A_n = -\frac{2a}{n\pi} (-1)^n (T_1 - T_0)$$

$$\frac{(T - T_0)}{(T_1 - T_0)} = \sum \frac{-2 \cancel{(T_1 - T_0)}}{n\pi} (-1)^n \cdot e^{-\left(\frac{n\pi}{a}\right)^2 Dt} \cdot \frac{1}{r} \sin \frac{n\pi}{a} r$$

Assignment 4, Problem 1



Ass 4, 2



Hollow conducting sphere, ^{radius a} bottom half grounded, top half set to V . Find potential ~~for~~ for $r > a$.

$\nabla^2 \Phi = 0$ and system has axisymmetry about z axis. Legendre Polynomials

$$\Phi(r, \theta) = \sum_{l=0}^{\infty} \left[A_l r^l + \frac{B_l}{r^{l+1}} \right] P_l(\mu) \quad (1)$$

notice that as $r \rightarrow \infty$ the $A_l r^l \rightarrow \infty$ NOT GOOD! We avoid this by setting $A_l = 0$

Now, find B by evaluating (1) at $r = a$ for known b.c. value

$$\Phi(a, \theta) = \sum_{l=0}^{\infty} \frac{B_l}{a^{l+1}} P_l(\mu) = \begin{cases} V & 0 \leq \mu \leq 1 \\ 0 & -1 \leq \mu < 0 \end{cases}$$

Use orthogonality

$$\int_{-1}^1 \sum_{l=0}^{\infty} \frac{B_l}{a^{l+1}} P_l(u) P_{l'}(u) du = \int_{-1}^1 v(u) P_{l'}(u) du$$

$$\frac{B_l}{a^{l+1}} \underbrace{\int_{-1}^1 P_l(u) P_l(u) du}_{\frac{2}{2l+1}} = v \int_0^1 P_l(u) du$$

$\frac{2}{2l+1}$ ← orthogonality condition

$$\frac{B_l}{a^{l+1}} \cdot \frac{2}{2l+1} = v \underbrace{\int_0^1 P_l(u) du}$$

integrate using step up/down properties ... see notes (my page 99)

$$= \frac{-v P_{l+1}(0)}{l}$$

$$B_l = \frac{(2l+1)a^{l+1}}{2} \cdot \frac{-v P_{l+1}(0)}{l}$$

$l=0$ is special case ↙


For $l=0$ $P_0(\mu) = 1$

$$\int_0^1 P_0(\mu) d\mu = \int_0^1 1 d\mu = 1$$

So

$$\frac{2B_0}{a} = V \int_0^1 P_0(\mu) d\mu = V$$

$$B_0 = \frac{aV}{2}$$

$$\Phi(r, \theta) = \frac{aV}{2r} \left[1 - \frac{2l+1}{l} \frac{a^l}{r^l} P_{l+1}(\cos\theta) \cdot P_l(\mu) \right]$$


watch out, must remember to multiply B_0 by $\frac{1}{r^{l+1}} = \frac{1}{r}$!

Ass 4, 2

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clear
close all
a = 1;
V = 1;

step = .05;
x = (-1:step:1)' * ones(1,length(-1:step:1));
y = x';

r = sqrt(x.^2+y.^2);

highs = find(r > 1);
r(highs) = NaN;

theta = atan2(y,x);
theta(highs) = NaN;

maxl = 20;

ttest = 0;
if (ttest == 0)

for ix = 1:length(x)
    for iy = 1:length(y)
        phi(ix,iy) = 0;
        rr = r(ix,iy);
        tt = theta(ix,iy);

        work = 0;
        for l = 1:maxl
            Plo = legendre(l+1,0);
            Plu = legendre(l,cos(tt));
            work = work + (2*l)/l * Plo(l) * (rr/a).^l * Plu(l);
            work = work + (2*l+1)/l * Plo(l) * (rr/a).^l * Plu(l);
        end
        phi(ix,iy) = V/2*(1-work);
    end
end

figure(1)
clf

imagesc(-1:step:1,-1:step:1,phi)
set(gca,'YDir','normal')
colorbar

figure(2)
clf
contour(-1:step:1,-1:step:1,phi)
set(gca,'YDir','normal')

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Inside



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colorbar
end

% Now do r > a!

step = .1;
mxval = 3;
x = (-mxval:step:mxval)' * ones(1,length(-mxval:step:mxval));
y = x';

r = sqrt(x.^2+y.^2);

lows = find(r < 1);
r(lows) = NaN;

theta = atan2(y,x);
theta(lows) = NaN;

maxl = 20;

for ix = 1:length(x)
    for iy = 1:length(y)
        phi(ix,iy) = 0;
        rr = r(ix,iy);
        tt = theta(ix,iy);

        work = 0;
        for l = 1:maxl
            Plo = legendre(l+1,0);
            Plu = legendre(l,cos(tt));
            work = work + (2*l)/l * Plo(1) * (a/rr).^(l+1) * (2*l + 1)/l * Plu(1);
            work = work + (2*l+1)/l * Plo(1) * a.^l * (1/rr).^(l) * Plu(1);
        end
        phi(ix,iy) = a*V/2/rr*(1-work);
        %phi(ix,iy) = a*V/2/rr*(1-work);
    end
end

figure(3)
clf

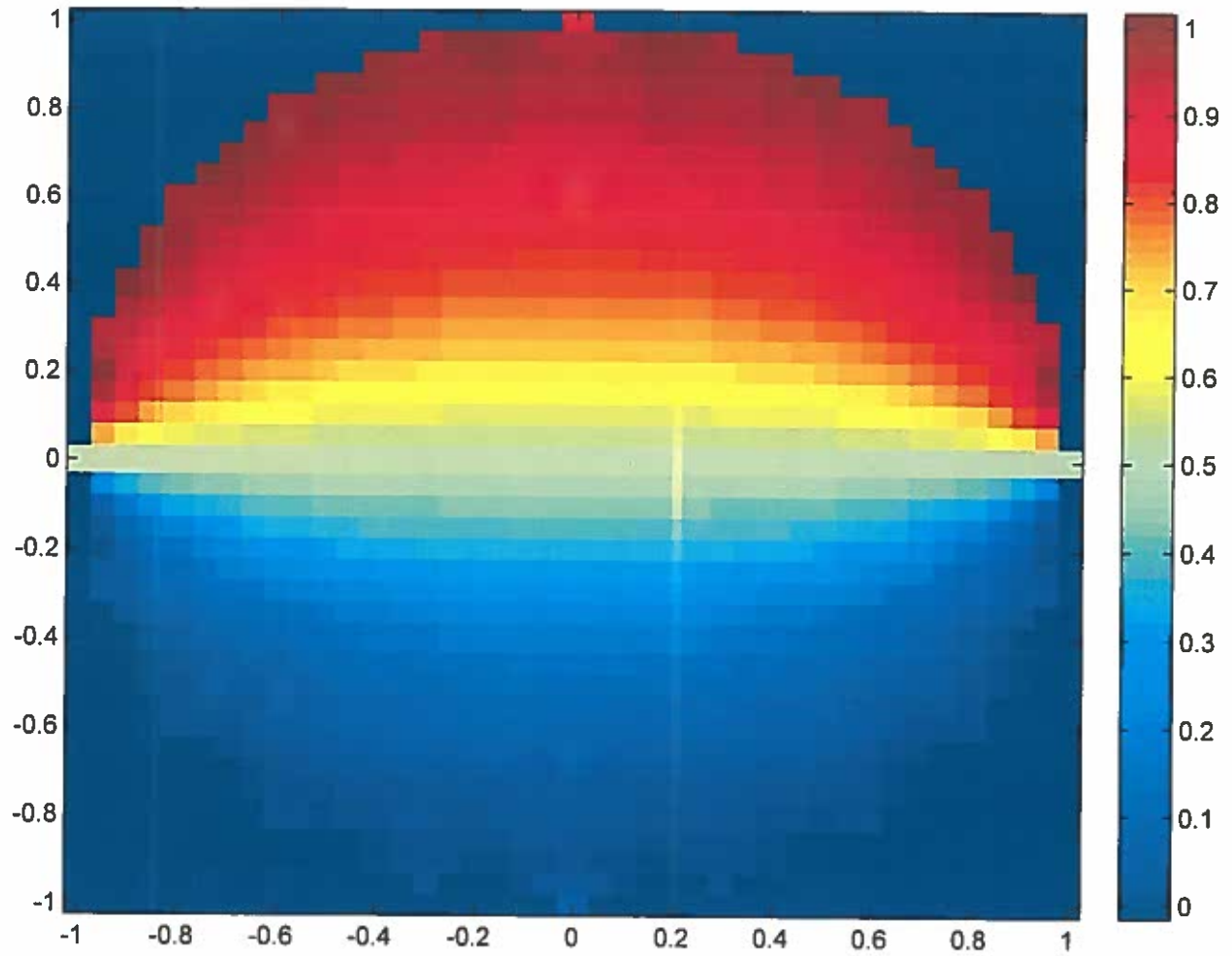
imagesc(-mxval:step:mxval,-mxval:step:mxval,phi)
set(gca,'YDir','normal')
colorbar

figure(4)
clf
contour(-mxval:step:mxval,-mxval:step:mxval,phi)
set(gca,'YDir','normal')
colorbar

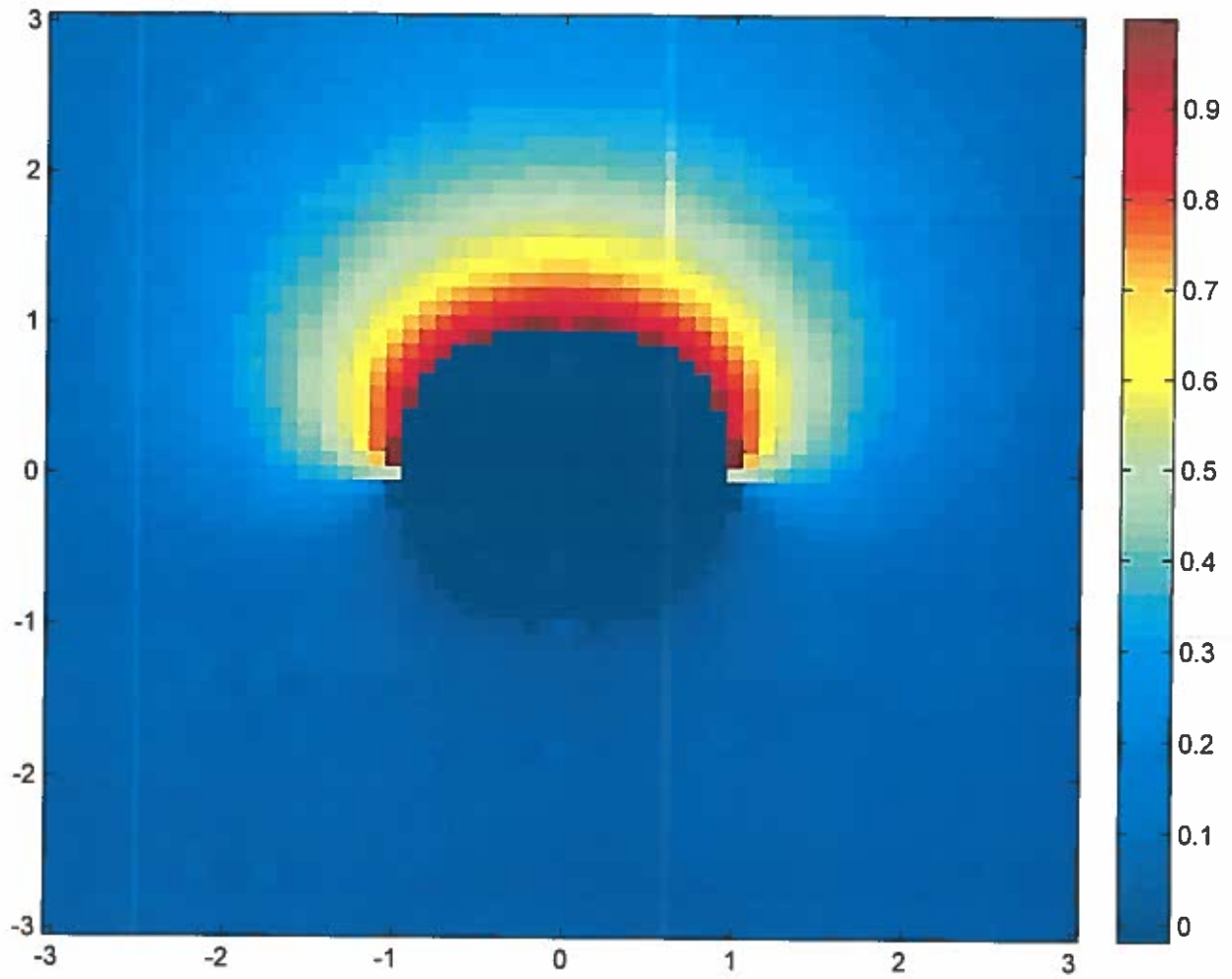
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Outside

Inside

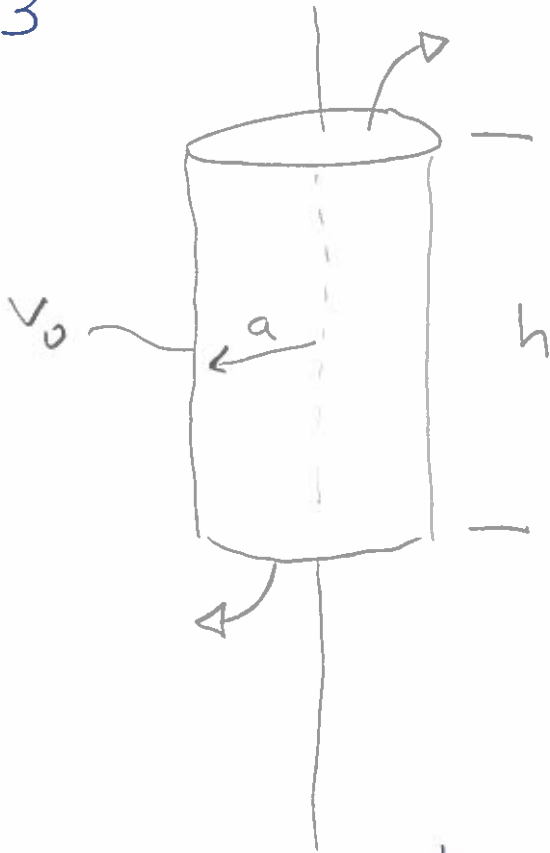


Outside



3

①



Want potential inside.

$$\nabla^2 \Phi = 0$$

This problem has axial symmetry so

$$w = e^{\pm im\phi}$$

has $m=0$.

General solution for $z = e^{ikz}$

need sin/cos form to match 0 at $z=0$ and $z=h$.

This form in z forces - sign in radial (ρ) eqn and you get solutions I_n and K_n .

K_n 's cannot be solutions because they are not finite at $\rho=0$. The only non-zero I_n

is I_0

General Form Solution:

$$\Phi(\rho, z) = \sum A_n \sin\left(\frac{n\pi z}{h}\right) I_0\left(\frac{n\pi}{h} \rho\right)$$

Set $\rho=a$ and find $\Phi(a, z) = V_0$ ∇

$$V_0 = \sum_n A_n \sin\left(\frac{n\pi z}{h}\right) I_0\left(\frac{n\pi}{h} a\right)$$

$$\frac{V_0}{I_0\left(\frac{n\pi}{h} a\right)} = \sum A_n \sin \frac{n\pi z}{h}$$

$$\int_0^h \frac{V_0}{I_0\left(\frac{n\pi}{h} a\right)} \cdot \sin\left(\frac{n\pi}{h} z\right) dz = A_n \int_0^h \sin^2\left(\frac{n\pi}{h} z\right) dz$$

$$\frac{V_0}{I_0\left(\frac{n\pi}{h} a\right)} \int_0^h \sin \frac{n\pi}{h} z dz = A_n \frac{h}{2}$$

$$\left[-\cos\left(\frac{n\pi}{h} z\right) \cdot \frac{h}{n\pi} \right]_0^h = -\cos n\pi - (-1) = \frac{h}{n\pi} (1 - \cos n\pi)$$

$$A_n \frac{h}{2} = \frac{V_0}{I_0\left(\frac{n\pi}{h} a\right)} \cdot \frac{h}{n\pi} (1 - \cos n\pi)$$

$\cos n\pi = -1$ n odd
 $\cos n\pi = 1$ n even

$$A_n = \frac{V_0}{I_0\left(\frac{n\pi}{h} a\right)} \cdot \frac{4}{n\pi}$$

3

$$\Phi(\rho, z) = \sum_{\substack{n \\ \text{odd}}} \frac{V_0 \cdot 4 \sin\left(\frac{n\pi}{h} z\right) \cdot I_0\left(\frac{n\pi}{h} \rho\right)}{I_0\left(\frac{n\pi}{h} a\right) n\pi}$$

$$\Phi(\rho, z) = \sum_n A_n \sin\left(\frac{n\pi z}{h}\right) I_0\left(\frac{n\pi}{h}\rho\right)$$

Now evaluate the potential at $\rho = a$:

$$V_0 = \sum_n A_n \sin\left(\frac{n\pi z}{h}\right) I_0\left(\frac{n\pi}{h}a\right)$$

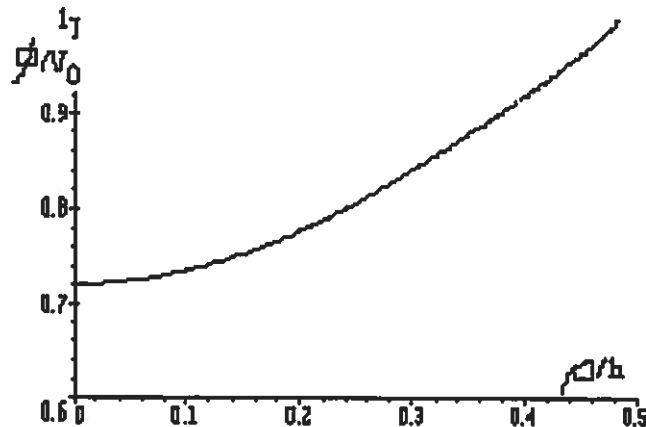
This is a Fourier sine series, and we find the coefficients in the usual way:

$$\begin{aligned} A_n \frac{h}{2} &= \frac{V_0}{I_0\left(\frac{n\pi}{h}a\right)} \int_0^h \sin\left(\frac{n\pi z}{h}\right) dz = \frac{V_0}{I_0\left(\frac{n\pi}{h}a\right)} \frac{h}{n\pi} (1 - \cos n\pi) \\ &= \frac{V_0}{I_0\left(\frac{n\pi}{h}a\right)} \frac{2h}{n\pi} \text{ for } n \text{ odd} \end{aligned}$$

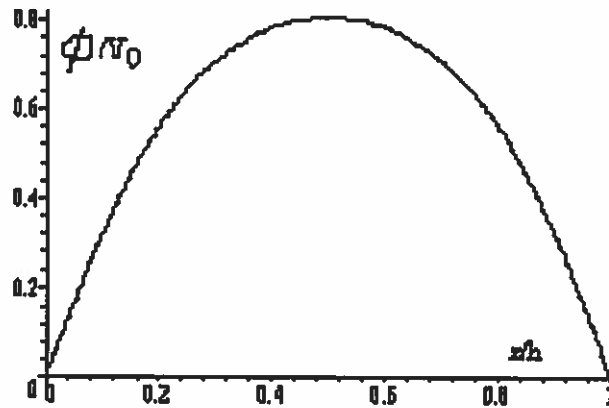
Thus

$$\Phi(\rho, z) = \frac{4V_0}{\pi} \sum_{n \text{ odd}} \frac{\sin\left(\frac{n\pi z}{h}\right)}{n} \frac{I_0\left(\frac{n\pi}{h}\rho\right)}{I_0\left(\frac{n\pi}{h}a\right)}$$

The first plot shows the first 9 terms in the expansion of $\Phi(\rho, h/2)/V_0$ for $a = h/2$.



The second plot shows $\Phi(a/2, z)/V_0$ with $a = h/2$.



43. A cylinder of height h and radius a is grounded except for its base at $z = 0$.

On the base the potential is $\frac{V_0}{\sqrt{1-\rho^2/a^2}}$. Find the potential inside the cylinder.

Again we have axisymmetry, and we must choose the eigenvalue to make the potential zero at $\rho = a$. To make the potential zero at $\rho = h$ we use the hyperbolic sine function.

$$\Phi(\rho, z) = \sum_n A_n \sinh\left(x_{0n} \frac{(h-z)}{a}\right) J_0\left(x_{0n} \frac{\rho}{a}\right)$$

where

$$\begin{aligned} A_n \sinh\left(\frac{h}{a} x_{0n}\right) \frac{a^2}{2} [J_0'(x_{0n})]^2 &= \int_0^a \frac{V_0}{\sqrt{1-\rho^2/a^2}} J_0\left(x_{0n} \frac{\rho}{a}\right) \rho d\rho \\ &= a^2 \int_0^1 \frac{V_0}{\sqrt{1-y^2}} J_0(x_{0n} y) y dy \\ &= V_0 a^2 \frac{\sin x_{0n}}{x_{0n}} \end{aligned}$$

We used GR 6.554#2 to evaluate the integral. . Thus

$$\sum_n 2V_0 \frac{\sin x_{0n}}{x_{0n}} \frac{\sinh\left(x_{0n} \frac{(h-z)}{a}\right)}{\sinh\left(\frac{h}{a} x_{0n}\right)} \frac{J_0\left(x_{0n} \frac{\rho}{a}\right)}{[J_0'(x_{0n})]^2}$$

44. (a) Use the series for $J_0(x)$ to show that its Laplace transform is