

P4820 Assignment II

Due, February 15, 2019

- 1) [10] Show that

$$\phi_n(x) = \frac{1 - \cos nx}{n\pi x^2}$$

is a delta sequence.

- 2) [5] Use delta sequences to show that the limit:

$$\lim_{l \rightarrow 0} \left[\frac{\delta(x) - \delta(x-l)}{l} \right]$$

exhibits the sifting property of $\delta'(x)$.

- 3) [5] A disk of radius α and mass M lies in the $x-y$ plane. Express the density in terms of delta functions in Cartesian coordinates, cylindrical coordinates, and spherical coordinates.
- 4) [10] A disk of charge with radius a and surface charge density $\sigma(r) = \sigma_0 r/a$ lies in the $x-y$ plane with the center at the origin. Find the volume charge density in a) cylindrical coordinates and b) spherical coordinates.
- 5) [10] A string of length L , with tension T and mass per unit length μ is hit simultaneously at $t = 0$ at the two points $x = L/3$ and $x = 2L/3$. The impulse delivered at each point is I . Find the equation that describes the subsequent displacement of the string. Plot the string position for times of $t = v * t/L = 1/10$ and $1/100$. Compare plots when you use the first 20 terms versus when you use the first 200 terms.

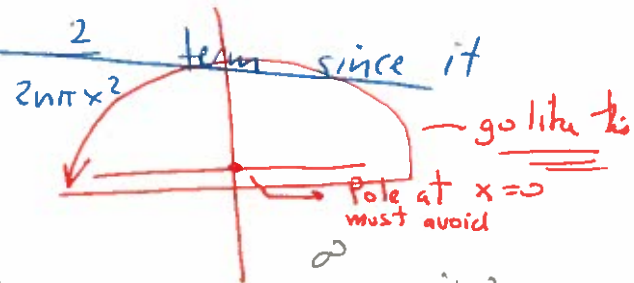
2)

$$\int_{-\infty}^{\infty} \left[\frac{1 - \cos nx}{n\pi x^2} \right] f(x) dx = ?$$

$$= \int_{-\infty}^{\infty} \frac{1}{n\pi x^2} f(x) dx - \int_{-\infty}^{\infty} \frac{e^{inx} + e^{-inx}}{2n\pi x^2} f(x) dx$$

$$\int_{-\infty}^{\infty} \frac{1}{n\pi x^2} \left(\frac{2 - e^{inx} - e^{-inx}}{2} \right) f(x) dx$$

This we will evaluate over the complex plane,
~~no need to fuss over the $\frac{2}{2n\pi x^2}$ term since it will go to zero at ∞~~



We consider

$$\underbrace{\int_{-\infty}^{\infty} \frac{1}{n\pi z^2} f(z) dz}_{\text{Residue} = 0!} - \int_{-\infty}^{\infty} \frac{e^{inz}}{2n\pi z^2} f(z) dz - \int_{-\infty}^{\infty} \frac{e^{-inz}}{2n\pi z^2} f(z) dz$$

close on top close on bottom

$$\frac{1}{2n\pi} \left(1 + inz - n^2 z^2 + \dots \right)$$

$$\frac{1}{z^2} + \frac{in}{z} - \dots$$

$$\text{Res} = \frac{-in}{2n\pi} \cdot \frac{+2\pi}{2n\pi} = 1 \quad \text{at } z=0$$

$$\boxed{= f(0)}$$

doesn't contain the pole at $z=0$

An alternative approach to the integral

$$\underbrace{\int_{-\infty}^{\infty} \frac{1}{n\pi z^2} f(z) dz}_{\text{Residue} = 0 \text{ no problem!}} - \int_{-\infty}^{\infty} \frac{e^{inz}}{2n\pi z^2} f(z) dz \quad \textcircled{1} - \int_{-\infty}^{\infty} \frac{e^{-inz}}{2n\pi z^2} f(z) dz \quad \textcircled{2}$$

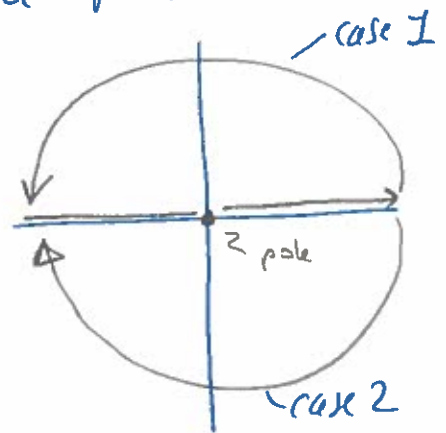
has pole at $x=0$! and that would be on my surface of integration if I keep the curve at $x=0$

In class we found that such a pole contributes half a residue!

$$e^{inz} = (1 + inz - n^2 z^2 + \dots)$$

$$\frac{e^{inz}}{2n\pi z^2} = \frac{1}{2n\pi z^2} + \frac{i}{2\pi z} - \dots$$

↑
residue



For integral $\textcircled{1}$, must close on top to keep integral bounded at ∞ .

$$\lim_{n \rightarrow \infty} \left[\int_{-\infty}^{\infty} \frac{e^{inz}}{2n\pi z^2} f(z) dz = -\frac{1}{2} \frac{i}{2\pi} \cdot 2\pi i = \frac{1}{2} f(0) \right]$$

$$\text{case } \textcircled{2} \quad \int_{-\infty}^{\infty} \frac{e^{-inz}}{2n\pi z^2} f(z) dz = -\frac{1}{2} \frac{i}{2\pi} \cdot 2\pi i = \frac{1}{2} f(0)$$

- sign from e^{-} but again from cw integration

Total integral

$$\int_{-\infty}^{\infty} \frac{1 - r \sin x}{h \pi x^2} f(x) dx = \frac{f(0)}{2} + \frac{f(0)}{2} = f(0)$$

Sifting property

$$3] \quad \lim_{l \rightarrow 0} \left[\frac{\delta(x) - \delta(x-l)}{l} \right]$$

Consider a delta sequence ϕ_n , $\lim_{n \rightarrow \infty} \phi_n \rightarrow \delta$

$$\int_{-\infty}^{\infty} \lim_{n \rightarrow \infty} \lim_{l \rightarrow 0} \left[\frac{\phi_n(x) - \phi_n(x-l)}{l} \right] f(x) dx$$

$$\lim_{n \rightarrow \infty} \lim_{l \rightarrow 0} \left[\frac{\int \phi_n(x) f(x) dx}{l} - \frac{\int \phi_n(x-l) f(x) dx}{l} \right]$$

$$= \lim_{l \rightarrow 0} \left[\frac{f(0) - f(+l)}{l} \right] = -f'(0)$$

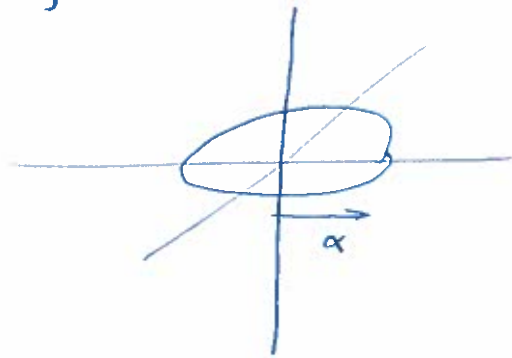
but, $\int \delta' f(x) dx = -f'(0)$

normally get $\frac{f(x+l) - f(x)}{l}$

$$\Rightarrow \boxed{\lim_{l \rightarrow 0} \frac{\delta(x) - \delta(x-l)}{l} = \delta'(x)}$$

4)

Disk mass M radius α in x, y plane
Find density!



i) Cartesian disk is at $z=0$

$$\rho(\vec{x}) = \frac{M}{\pi\alpha^2} \delta(z) \quad \text{for } \sqrt{x^2+y^2} < \alpha$$

and 0 otherwise:

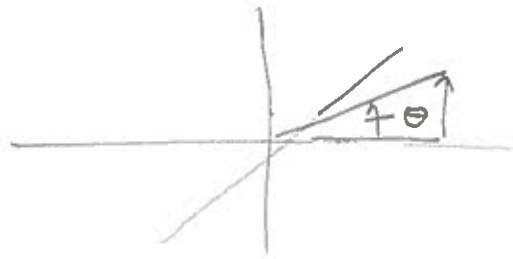
$$\rho(\vec{x}) = \frac{M}{\pi\alpha^2} \delta(z) \Theta(\alpha - \sqrt{x^2+y^2})$$

↑
step function

ii) Cylindrical coordinates
easy modification to cartesian, $\sqrt{x^2+y^2} = r$

$$\rho(\vec{x}) = \frac{M}{\pi\alpha^2} \delta(z) \Theta(\alpha - r)$$

iii) Spherical



$$\rho(\vec{x}) = \frac{M}{\pi d^2} \delta(r \cos \theta) \Theta(\alpha - r) \quad z \Rightarrow r \cos \theta$$

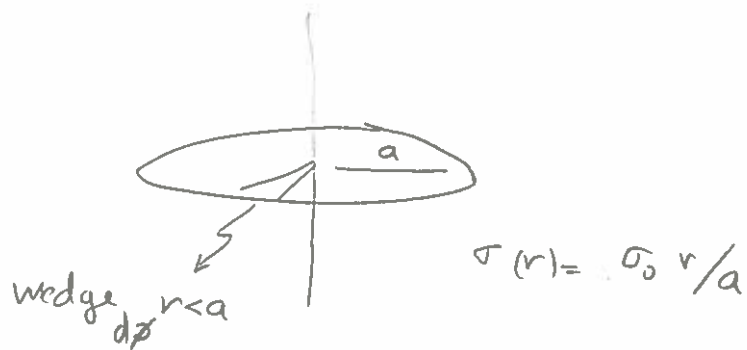
$$= \frac{M}{\pi d^2 r} \delta(\cos \theta) \Theta(\alpha - r)$$

↓
has a zero at $\theta = \frac{\pi}{2}$

$$= \frac{M}{\pi d^2 r} \frac{\delta(\theta - \frac{\pi}{2})}{|-\sin(\frac{\pi}{2})|} \Theta(\alpha - r)$$

$$\rho(\vec{x}) = \frac{M}{\pi d^2 r} \delta(\theta - \frac{\pi}{2}) \Theta(\alpha - r)$$

5)



i) In cylindrical coordinates

$$\rho(\vec{x}) = f(r, z) \sigma_0 \Theta(a-r) \delta(z)$$

integrate over a wedge of a cylindrical shell $r < a$

$$dq = \sigma_0 \frac{r^2}{a} dr d\phi = \int_{-\infty}^{\infty} f(r, z) \sigma_0 \Theta(a-r) \delta(z) r dr d\phi dz$$

$$\sigma_0 \frac{r^2}{a} dr d\phi = f(r, 0) \sigma_0 r dr d\phi$$

$$f(r, 0) = \frac{r}{a}$$

$$\rho(\vec{x}) = \frac{\sigma_0 r}{a} \Theta(a-r) \delta(z)$$

ii) In polar coordinates now integrate over an orange wedge slice

$$\rho(\vec{x}) = f(r, \theta, \phi) \sigma_0 \Theta(a-r) \delta(r \cos \theta) \quad \left. \begin{array}{l} \text{absorbs } r \\ \text{into unknown } f \end{array} \right\}$$

$$= f(r, \theta) \sigma_0 \Theta(a-r) \delta(\theta - \pi/2)$$

$$\frac{\sigma_0 r^2}{a} dr d\phi = \int_0^\pi f(r, \theta) \sigma_0 \Theta(a-r) \delta(\theta - \pi/2) r^2 dr d\phi \sin \theta d\theta$$

$$= f(r, \pi/2) \sigma_0 \int r^2 dr d\phi \cdot 1$$

$$f(r, \pi/2) = \frac{1}{a}$$

$$\rho(\vec{x}) = \frac{\sigma_0}{a} \Theta(a-r) \delta(\theta - \pi/2)$$

again, by just converting variables
 $z = r \cos \theta$!

$$\rho(\vec{x}) = \frac{\sigma_0 r}{a} \Theta(a-r) \delta(z)$$

↪ to ~~cylindrical~~ spherical



$$\rho(\vec{x}) = \frac{\sigma_0 r}{a} \Theta(a-r) \delta(r \cos \theta)$$

$$= \frac{\sigma_0 r}{a r} \Theta(a-r) \delta(\cos \theta)$$

$$= \frac{\sigma_0}{a} \Theta(a-r) \frac{\delta(\theta - \pi/2)}{|-\sin \theta|}$$

$$= \frac{\sigma_0}{a} \Theta(a-r) \delta(\theta - \pi/2)$$

again

6) A string length L , tension T and mass density μ .

The initial condition we have at $t=0$

$$\left. \frac{\partial y}{\partial t} \right|_{t=0} = \frac{F}{\mu} \left[\delta\left(x - \frac{L}{3}\right) + \delta\left(x - \frac{2L}{3}\right) \right]$$

also, $y(x, 0) = 0$

and of course $y(0, t) = y(L, t) = 0$

we seek a solution of the form:

$$y(x, t) = \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{L} \sin \frac{n\pi v t}{L}$$

$$v = \sqrt{\frac{T}{\mu}}$$

gotta work from $y'(x, t) =$

$$= \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{L} \cdot \cos \frac{n\pi v t}{L} \cdot \frac{n\pi v}{L} \Big|_{t=0}$$

$$\frac{F}{\mu} \left[\delta\left(x - \frac{L}{3}\right) + \delta\left(x - \frac{2L}{3}\right) \right] = \sum_{n=1}^{\infty} a_n \frac{n\pi v}{L} \cdot \sin \frac{n\pi x}{L}$$

Integrate to find a_n

$$\frac{2}{L} \frac{I}{\mu} \int_{-\infty}^{\infty} \sin \frac{n\pi x}{L} \left[\delta\left(x - \frac{L}{3}\right) + \delta\left(x - \frac{2L}{3}\right) \right]$$

$$= \frac{2}{L} \frac{I}{\mu} \left(\sin \frac{n\pi}{3} + \sin \frac{2n\pi}{3} \right) = a_n \cdot \frac{n\pi v}{L}$$

$$a_n = \frac{2}{n\pi v} \cdot \frac{I}{\mu} \left(\sin \frac{n\pi}{3} + \sin \frac{2n\pi}{3} \right)$$

$$y(x, t) = \sum_{n=1}^{\infty} \frac{2I}{n\pi v \mu} \left(\sin \frac{n\pi}{3} + \sin \frac{2n\pi}{3} \right) \sin \frac{n\pi x}{L} \sin \frac{n\pi v t}{L}$$

Does not really converge that fast
 ... $\frac{1}{n}$ term!

